

# M285K - course #15

# IV. QUANTITATIVE THEORY:

## FLUCTUATIONS.

$$- \nabla \cdot a(\dot{\varepsilon}) \nabla u_\varepsilon = \nabla \cdot f$$

3 questions: \* macro observable  $\int g \cdot \nabla u_\varepsilon \xrightarrow{\text{a.s.}} \int g \cdot \nabla \bar{u}$   
in terms of some  
homog eqn.

\*  $a(\dot{\varepsilon})$  oscillating on scale  $O(\varepsilon)$   
 $\Rightarrow \nabla u_\varepsilon$  should also on scale  $O(\varepsilon)$ .

Intrinsic description: 2-scale expansion

$$\nabla u_\varepsilon - (\nabla a(\dot{\varepsilon}) + e_i) \nabla_i \bar{u} = O(\varepsilon) \quad \underline{d \geq 2}$$

in  $L^p(\mathbb{R}^d, \mathcal{L}^q(\Omega))$   
 $1 < p, q < \infty$ .

+ get a rate:  $\left| \int g \cdot \nu_{u_\varepsilon} - \int g \cdot \bar{\nu} \right|$  in  $L^q(\Omega), q < \infty$

$$\leq O(\varepsilon) + \underbrace{\left| \int g \cdot \nu_{\varphi_i(\bar{i})} \bar{\nu} \right|}_{O(\varepsilon^{d/2}) \text{ in } L^q(\Omega)!}$$

(of CLT holds for  $\nu_\varphi$ !)

⊗ a random  $\Rightarrow \nu_{u_\varepsilon}$  random

$$\varepsilon^{-d/2} \int g \cdot (\nu_{\bar{i}} - \mathbb{E} \nu_{\bar{i}}) \xrightarrow[\text{(CLT)}]{\text{law}} \text{Gaussian } \int g \cdot \bar{\nu} \xi$$

white noise.

↳ expect:  $\varepsilon^{-d/2} \int g \cdot (\nu_{u_\varepsilon} - \mathbb{E} \nu_{u_\varepsilon}) \xrightarrow{\text{law}} \text{Gaussian?}$   
 should not be white noise!

Good get some intrinsic description!

$\frac{1}{\sqrt{\epsilon}}$

2-side expansion? -

Observation.

$\left\{ \begin{array}{l} \epsilon^{-d/2} \int g. (v_{u_\epsilon} - \mathbb{E} v_{u_\epsilon}) \xrightarrow{\text{low}} \text{Gaussian} \end{array} \right.$

$\left\{ \begin{array}{l} \epsilon^{-d/2} \int g. \nabla \varphi_i(i) \nabla_i \bar{u} \xrightarrow{\text{low}} \text{Gaussian} \end{array} \right.$

BUT the limits are different!

→ 2-side expansion is NOT accurate in fluctuation regime ---

Question: [do we still have an intrinsic description in terms of connectors?

Key quantity: homogenization commutator

commutator between  
- large-scale average

$$\left[ \begin{array}{c} \mathcal{H}' \\ \mathcal{U}_\varepsilon \end{array} \right] [\psi_{u_\varepsilon}] := a(\bar{z}) \psi_{u_\varepsilon} - \bar{a} \psi_{u_\varepsilon} \quad \left. \begin{array}{l} \\ \end{array} \right\} \text{field-flux version}$$

Recall that this is a natural quantity:

$$\left[ \begin{array}{l} \left[ \begin{array}{c} \mathcal{H}' \\ \mathcal{U}_\varepsilon \end{array} \right] [\psi_{u_\varepsilon}] \rightarrow 0 \text{ in } L^2(\mathbb{R}^d)^d \\ \Leftrightarrow \psi_{u_\varepsilon} \rightarrow \bar{\psi} \text{ \& } a(\bar{z}) \psi_{u_\varepsilon} \rightarrow \bar{a} \bar{\psi} \text{ in } L^2(\mathbb{R}^d)^d \end{array} \right\} \begin{array}{l} d \text{ relation} \\ 2d \text{ relation.} \end{array}$$

3 key principles (to be stated later):

①  $\begin{pmatrix} \psi_{u_\varepsilon} \\ a(\bar{z}) \psi_{u_\varepsilon} \end{pmatrix} = \text{deterministic projection of } \left[ \begin{array}{c} \mathcal{H}' \\ \mathcal{U}_\varepsilon \end{array} \right] [\psi_{u_\varepsilon}]$

↳ all fluctuations reduce to fluctuations of  $\left[ \begin{array}{c} \mathcal{H}' \\ \mathcal{U}_\varepsilon \end{array} \right] [\psi_{u_\varepsilon}]$ .

② 2-side expansion of commutator:

$$\boxed{H}_\varepsilon[\psi_\varepsilon] \sim \boxed{H}_\varepsilon \left[ \underbrace{(\nabla\psi_i(\frac{i}{\varepsilon}) + e_i)}_{\boxed{H}_i^0: \text{"standard commutator"}} \psi_i \bar{u} \right] + O(\varepsilon) \quad (d \geq 2)$$

in fluctuation  
scaling!

(res:  $\psi_\varepsilon \sim (\nabla\psi_i(\frac{i}{\varepsilon}) + e_i) \psi_i \bar{u} + O(\varepsilon)$  in fluct scaling)

③  $\boxed{H}_\varepsilon^0$  is "almost" local w.r.t.  $\varepsilon$ .

$\Rightarrow \boxed{H}_\varepsilon^0 \sim$  white noise on large scales —  $\mathcal{M}$

Remark: role of commutator for fluctuations = improved locality.

$$-\nabla \cdot a \nabla u = \nabla \cdot f \quad (\varepsilon = 1)$$

⊗  $\nabla u$ : critical load wrt  $a$ .

Mollifier:  $-\nabla \cdot a \nabla D_z u = \nabla \cdot D_a \nabla u$

$$a(x) = a_0(\theta(x))$$

$$D_z \nabla u(x) = \underbrace{\nabla G(x, z)}_{\sim |x-z|^{-d}} a'_0(\theta(z)) \nabla u(z)$$

⊗  $\square[\nabla u]$ : exactly load wrt  $a$  + small error.

Lemma. 
$$\begin{aligned} \textcircled{1} \square[\nabla u]_j &= (\nabla \varphi_j^* + e_j) \cdot \overset{\delta}{D_a} \nabla u \\ &\quad - \underbrace{\nabla \cdot \left( (\varphi_j^* a + \sigma_j^*) \nabla D u \right)}_{\substack{\delta \\ \text{small} \\ \text{or large scales!}}} + \varphi_j^* \underbrace{D_a \nabla u}_{\delta} \end{aligned}$$

(1)  $\square[\nabla u]_j$

$\delta$  small or large

↳ open & perturbation.

$$D \mathbb{H}_\varepsilon [v_{u_j}] = (v_{\psi_j}^* |_{\dot{\varepsilon}} + e_j) \cdot D a(\dot{\varepsilon}) v_{u_j} - \nabla \cdot \left( \varepsilon (v_{\psi_j}^* + \sigma_j^*) |_{\dot{\varepsilon}} \nabla v_{u_j} + \varepsilon v_{\psi_j}^* |_{\dot{\varepsilon}} D a(\dot{\varepsilon}) v_{u_j} \right)$$

$\mathcal{O}(\varepsilon)$ , critical locality

Proof.  $\mathbb{H}[v_u] = a v_u - \bar{a} v_u$

$$\left\{ \begin{aligned} -\nabla \cdot a^* (v_{\psi_j}^* + e_j) &= 0 \end{aligned} \right.$$

$$\left\{ \begin{aligned} \nabla \cdot \sigma_j^* &= a^* (v_{\psi_j}^* + e_j) - \bar{a}^* e_j. \end{aligned} \right.$$

$$\rightarrow D \mathbb{H}_\varepsilon [v_{u_j}] = e_j \cdot D a v_u + \underbrace{e_j \cdot (a - \bar{a}) \nabla v_u}_{(a^* - \bar{a}^*) e_j \cdot \nabla v_u}$$

Decompose:  $(a^* - \bar{a}^*)e_j = \underbrace{-a^* \nabla \psi_j}_{\text{grad like}} + \underbrace{(a^* (\nabla \psi_j + e_j) - a^* e_j)}_{\nabla \cdot \sigma_j^*}$

"a-Helmholtz decomposition" of  $a^* - \bar{a}^*$ !

$$\begin{aligned} \int_{\Omega} (a^* - \bar{a}^*)e_j \cdot \nabla Du &= - \int_{\Omega} \nabla \psi_j^* \cdot a^* \nabla Du + \int_{\Omega} (\nabla \cdot \sigma_j^*) \cdot \nabla Du \\ &= - \int_{\Omega} \nabla \cdot (\psi_j^* a^* \nabla Du) + \int_{\Omega} \psi_j^* \nabla \cdot a^* \nabla Du \\ &\quad - \int_{\Omega} \nabla \cdot (\sigma_j^* \nabla Du) \end{aligned}$$

where  $\int_{\Omega} (\nabla \cdot \sigma_j^*) \cdot \nabla Du = (\nabla_e \sigma_{jkl}^*) (\nabla_a Du)$

$$= \int_{\Omega} \sigma_{ihl}^* \nabla_h Du - \int_{\Omega} \sigma_{ihl}^* \nabla_{al}^2 Du$$

$\nabla \cdot (\sigma_{jlk})$  show sym

$$= -\nabla \cdot (\sigma_j^* \nabla Du)$$

$$\rightarrow (\bar{a} - \bar{a}^*) e_j \cdot \nabla Du = -\nabla \cdot (\varphi_j^* a + \sigma_j^*) \nabla Du$$

$$\begin{aligned}
 &+ \varphi_j^* \underbrace{\nabla \cdot a \nabla Du}_{-\nabla \cdot Da \nabla Du} \\
 &- \nabla \cdot (\varphi_j^* Da \nabla Du) \\
 &+ \nabla \varphi_j^* \cdot Da \nabla Du
 \end{aligned}$$

### IV.1 Pathwise structure of fluctuations

+ 1 + d.p.t. in time of 25 node economy -

= intrinsic description of functions in any of ...

Lemma (Principle 1: reduction to commutator).

$$\int_{\mathbb{R}^d} g \cdot (\nabla u_\varepsilon - \nabla \bar{u}) = \int_{\mathbb{R}^d} \bar{\mathcal{H}}^*(g) \cdot \overline{\mathcal{C}}_\varepsilon[\nabla u_\varepsilon]$$

$$\int_{\mathbb{R}^d} g \cdot (\mathfrak{a}(\varepsilon) \nabla u_\varepsilon - \bar{\mathfrak{a}} \nabla \bar{u}) = \int_{\mathbb{R}^d} \bar{\mathcal{L}}^*(g) \cdot \overline{\mathcal{C}}_\varepsilon[\nabla u_\varepsilon]$$

where

- $\bar{\mathcal{H}}(h) = \nabla(-\nabla \cdot \bar{\mathfrak{a}} \nabla)^{-1} \nabla \cdot h$   $\bar{\mathfrak{a}}$ -Helmholtz projection (grad-like)
- $\bar{\mathcal{L}}(h) = h + \bar{\mathfrak{a}} \nabla(-\nabla \cdot \bar{\mathfrak{a}} \nabla)^{-1} \nabla \cdot h$   $\bar{\mathfrak{a}}$ -Leray projection (div-free)

Prop 1  $1 - \nabla \cdot \bar{\mathfrak{a}}^* \nabla \bar{\mathfrak{a}} = \nabla \cdot \bar{\mathfrak{g}}$



derog proj  $\Delta_0 f = f - \frac{1}{V} \nabla \cdot f$  )  $V \cdot \Delta_0 f = \nabla \cdot f - \frac{1}{V} \nabla \cdot \nabla \cdot f$

div free = 0

$L^2$   
 $\downarrow$   
 $f = \underbrace{\Delta_0 f}_{\text{grad}} + \underbrace{\Delta_0 f}_{\text{div free}}$