

Recap: Effective viscosity $e \cdot \bar{B} e = E |D\varphi_e + e|^2$
 $= |e|^2 + E |D\varphi_e|^2 \gg |e|^2$
 (larger than solvent viscosity
 as particles hinder fluid flow)

- defined in terms of connector $\varphi \Rightarrow$ depends on the law of microstructure $\{I_n\}_n$



... and non-explicit!

- needs to be coupled with microstructure dynamics

\hookrightarrow fully open problem $\left\{ \begin{array}{l} \text{fast micro dynamics} \\ \text{infinite particle system} \dots \end{array} \right.$

- simplification: dilute regime \Rightarrow only reduced info on microstructure matters
 $\&$ get (semi)explicit perturbative description of \bar{B}

Dilute means: volume fraction $\lambda := E \mathbb{1}_y = \lim_{R \uparrow \infty} \frac{f}{B_R} \mathbb{1}_y = \lim_{R \uparrow \infty} \frac{|y \cap B_R|}{|B_R|}$ e.s.
 $\ll 1$.

II Dilute expansions.

Long history: 19th century, macro effect of impurities in physical systems
 (Clausius, Monetti, Maxwell, Rayleigh, Faraday ...)

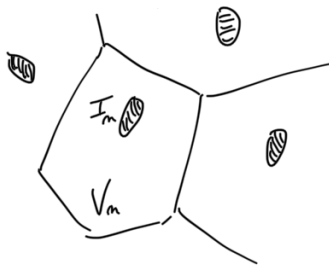
* Einstein 1905: $\bar{B} = \underbrace{Id}_{\text{pure solvent}} + \underbrace{\frac{5}{2} \lambda Id}_{O(\lambda) \text{ contribution from } \text{1-particle effects}} + o(\lambda)$ (3D, Stokes spheres)

* next orders: nnnnn 2-particle effects, etc "cluster expansion"
 BUT divergent! (long-range interactions)
 e.g. Burgers '40s...

\hookrightarrow solved to 2nd order: Batchelor-Green '72 - renormalization.

II.1) Einstein's formula.

Heuristics:



dilute regime \rightarrow particles as isolated...

$\psi_e|_{V_m} \approx \psi_e^m$ solution of single-particle problem

$$\hookrightarrow \psi_e \approx \sum_n \mathbb{1}_{V_n} \psi_e^m$$

$$\hookrightarrow e \cdot \bar{B} e \approx \text{Id} + \mathbb{E} \sum_n \mathbb{1}_{V_n} |\psi_e^m|^2 \approx \text{Id} + \mathbb{E} \sum_n \frac{\mathbb{1}_{I_n}}{|I_n|} \int_{\mathbb{R}^d} |\psi_e^m|^2$$

Theorem.

For simplicity: $\left\{ \begin{array}{l} - \text{scalar setting} \\ - \text{assume (H_{unif})} : \text{dist}(I_m, I_n) \geq \delta > 0 \forall m \neq n \text{ a.s.} \end{array} \right.$

$$\text{then } \bar{B} = \text{Id} + \bar{B}^{(n)} + O(\lambda^2 \log(2 + \frac{1}{\lambda}))$$

\uparrow think λ^2 for Poincaré (hardcore) - see later!

$$\text{where } e \cdot \bar{B}^{(n)} e = \mathbb{E} \sum_n \frac{\mathbb{1}_{I_n}}{|I_n|} \int_{\mathbb{R}^d} |\psi_e^m|^2, \quad \left\{ \begin{array}{l} - \Delta \psi_e^m = 0, I_n^c \\ \psi_e^m|_{I_n} = 0, I_n \\ \int_{\partial I_n} \partial_n \psi_e^m = 0 \end{array} \right.$$

$$= O(\lambda)$$

In particular:

* if $\{X_n\}_n \perp \text{iid } \{R_n\} \sim R_0, I_n = I(X_n, R_n),$

$$\text{then } e \cdot \bar{B}^{(n)} e = \frac{1}{|I|} \mathbb{E}_{R_0} \int_{\mathbb{R}^d} |\psi_e^0|^2, \quad \left\{ \begin{array}{l} - \Delta \psi_e^0 = 0, I(Q, R_0)^c \\ \psi_e^0|_{I(Q, R_0)} = 0, I(Q, R_0) \\ \int_{\partial I(Q, R_0)} \partial_n \psi_e^0 = 0 \end{array} \right.$$

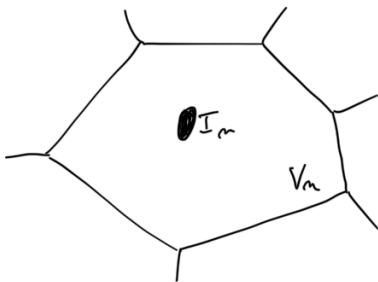
* if $I_n = B(x_n),$ then $\bar{B}^{(n)} = \text{Id}$

(Stokes $\rightarrow \lambda \frac{d+2}{2} \text{Id} : \text{Einstein formula}$)

Proof: can do PDE proof, but variational proof is most natural.

$$\text{Start from: } e \cdot \bar{B} e = \mathbb{E} |\psi_{e+e}|^2 = \inf \left\{ \mathbb{E} |\psi_{y+c}|^2 : \psi_y \text{ stat, } \psi_{y+c}|_{I_n} = 0, \mathbb{E} \psi_y = 0 \right\}$$

$$= |e|^2 + \mathbb{E} |\psi_y|^2$$



Voronoi cells: $V_m = \{x \in \mathbb{R}^d : |x - x_m| \leq |x - x_n| \forall n \neq m\}$

Single-particle problem:

$$\text{① whole-space: } \left\{ \begin{array}{l} - \Delta \psi_e^m = 0, I_n^c \\ \psi_e^m|_{I_n} = 0 \\ \int_{\partial I_n} \partial_n \psi_e^m = 0, \psi_e^m \in H^1(\mathbb{R}^d) \end{array} \right.$$

$$\text{i.e. } \psi_e^m = \text{argmin} \left\{ \int_{\mathbb{R}^d} |\psi_y|^2 : y \in H^1(\mathbb{R}^d), \psi_y|_{I_n} = 0 \right\}$$

2) Dirichlet :

$$\varphi_{e,D}^m = \operatorname{argmin} \left[\int_{V_m} |\nabla \varphi|^2 : \varphi \in H_0^1(V_m), \nabla \varphi + e|_{\Gamma_m} = 0 \right]$$

3) Neumann :

$$\varphi_{e,N}^m = \operatorname{argmin} \left\{ \int_{V_m} |\nabla \varphi|^2 : \varphi \in H^1(V_m), \nabla \varphi + e|_{\Gamma_m} = 0, \int_{V_m} \varphi = 0 \right\}$$

Step 1 Energy comparison:

$$\begin{cases} e \cdot \bar{B} e \leq |e|^2 + E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \int_{V_n} |\nabla \varphi_{e,D}^n|^2 \\ e \cdot \bar{B} e \geq |e|^2 + E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \int_{V_n} |\nabla \varphi_{e,N}^n|^2 \end{cases}$$

Exercise: averaging over Voronoi cells $\rightarrow \left[\forall \xi \text{ stationary}, E|\xi| < \infty, \right.$
 $\left. \text{we have } E\xi = E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \int_{V_n} \xi \right]$

⊆ Consider $\varphi_{e,D} = \sum_n \varphi_{e,D}^n$ stationary

$$\begin{cases} \text{Get } E \nabla \varphi_{e,D} = E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \underbrace{\int_{V_n} \nabla \varphi_{e,D}^n}_{=0} = 0! \\ E |\nabla \varphi_{e,D}|^2 = E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \underbrace{\int_{V_n} |\nabla \varphi_{e,D}^n|^2}_{\leq 1} \leq 1 \end{cases}$$

$$\begin{aligned} \Rightarrow \text{by energy comparison: } e \cdot \bar{B} e &\leq |e|^2 + E |\nabla \varphi_{e,D}|^2 \\ &= |e|^2 + E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \int_{V_n} |\nabla \varphi_{e,D}^n|^2. \end{aligned}$$

$$\begin{aligned} \text{⊇ } e \cdot \bar{B} e &= |e|^2 + \sum_n E \frac{|\Gamma_n|}{|\Gamma_n|} \underbrace{\int_{V_n} |\nabla \varphi_{e,N}^n|^2}_{\geq \int_{V_n} |\nabla \varphi_{e,D}^n|^2} \\ &\geq \int_{V_n} |\nabla \varphi_{e,D}^n|^2! \quad - \end{aligned}$$

$$\text{Step 2 } \left(\int_{V_m} |\nabla \varphi_{e,N}^m|^2 \right) \leq \int_{V_m} |\nabla \varphi_{e,D}^m|^2 \leq \left(\int_{\mathbb{R}^d} |\nabla \varphi_{e,N}^m|^2 \right) \leq \int_{\mathbb{R}^d} |\nabla \varphi_{e,D}^m|^2 = \left(\int_{V_m} |\nabla \varphi_{e,D}^m|^2 \right)$$

$$\hookrightarrow \left| e \cdot \bar{B} e - |e|^2 - E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \int_{\mathbb{R}^d} |\nabla \varphi_{e,N}^n|^2 \right| \leq E \sum_n \frac{|\Gamma_n|}{|\Gamma_n|} \left[\int_{V_m} |\nabla \varphi_{e,D}^n|^2 - |\nabla \varphi_{e,N}^n|^2 \right]$$

$$\text{Step 3 } \left[\int_{V_m} |\nabla \varphi_{e,D}^n|^2 - |\nabla \varphi_{e,N}^n|^2 \right] \leq (\rho_m)^{-d},$$

by elliptic regularity!

$$\begin{aligned} \rho_m &= \max \{ r : B_r(x_m) \subset V_m \} \\ &= \frac{1}{2} \min_{m:m+n} |x_m - x_n| \end{aligned}$$

$$\begin{aligned}
\bullet \int_{V_m} |\nabla \varphi_{eD}^m|^2 - |\nabla \varphi_{eN}^m|^2 &= \int_{V_m \setminus I_m} \operatorname{div}(\varphi_{eD}^m \nabla \varphi_{eD}^m) - \operatorname{div}(\varphi_{eN}^m \nabla \varphi_{eN}^m) \\
&\stackrel{SC}{=} \int_{\partial I_m} e \cdot (x - x_m) \partial_n (\varphi_{eD}^m - \varphi_{eN}^m) \\
&\leq \left(\int_{I_m \setminus B} |\nabla(\varphi_{eD}^m - \varphi_{eN}^m)|^2 \right)^{1/2} \quad \text{by Exercise 1 (trace)} \\
&\leq \rho_m^{-d/2} \left(\int_{V_m} |\nabla(\varphi_{eD}^m - \varphi_{eN}^m)|^2 \right)^{1/2} \quad \text{by Exercise 2 (MVP)}
\end{aligned}$$

$$\begin{aligned}
\bullet \int_{V_m} |\nabla(\varphi_{eD}^m - \varphi_{eN}^m)|^2 &= \int_{V_m \setminus I_m} \operatorname{div}(\varphi_{eD}^m \nabla \varphi_{eD}^m) + \operatorname{div}(\varphi_{eN}^m \nabla \varphi_{eN}^m) - 2 \operatorname{div}(\varphi_{eD}^m \nabla \varphi_{eN}^m) \\
&= \int_{\partial I_m} e \cdot (x - x_m) \partial_n (\varphi_{eD}^m - \varphi_{eN}^m) \\
\hookrightarrow \int_{V_m} |\nabla(\varphi_{eD}^m - \varphi_{eN}^m)|^2 &= \int_{V_m} |\nabla \varphi_{eD}^m|^2 - |\nabla \varphi_{eN}^m|^2 \\
&\quad \& \text{claim follows:}
\end{aligned}$$

Exercise 1: trace estimate

$$\begin{cases} -\Delta \varphi = 0, & \partial B \setminus B \\ \nabla \varphi = 0, & B \end{cases}
\Rightarrow \int_{\partial B} |\nabla \varphi|^2 \leq \int_{\partial B \setminus B} |\nabla \varphi|^2 :$$

indeed: $\int_{\partial B} |\nabla \varphi|^2 \lesssim \int_{\partial B \setminus B} |\langle \nabla \varphi \rangle|^2$ by usual trace estimate $H^{1/2}(\partial B \setminus B) \hookrightarrow L^2(\partial B)$

$$\lesssim \int_{\partial B \setminus B} |\nabla \varphi|^2 + \underbrace{\int_{\partial B} |\langle \nabla \varphi \rangle \cdot (x - x_m)|^2}_{\substack{\text{can subtract} \\ \text{any constant}}} \quad \text{by elliptic boundary regularity}$$

= 0 for suitable κ as $\varphi|_B$ is const

Exercise 2: mean-value property (MVP)

recall for harmonic fcts: $\begin{cases} -\Delta \varphi = 0, & B_R \\ \Rightarrow |\nabla \varphi(0)|^2 \lesssim R^{-d} \int_{B_R} |\nabla \varphi|^2 \end{cases}$

deduce that we have similarly $\begin{cases} -\Delta \varphi = 0, & B_R \setminus B, \quad R \geq 2 \\ \nabla \varphi = 0, & B \\ \int_{B_R \setminus B} \varphi = 0 \end{cases} \Rightarrow \int_{B_R} |\nabla \varphi|^2 \lesssim R^{-d} \int_{B_R} |\nabla \varphi|^2$

indeed consider $\begin{cases} -\Delta \hat{\varphi} = 0, & B_R \quad \text{harmonic approximation} \\ \hat{\varphi} = \varphi, & \partial B_R \end{cases}$

standard MVP: $\int_{\partial B} |\nabla \hat{\varphi}|^2 \leq R^{-d} \int_{B_R} |\nabla \hat{\varphi}|^2$

$$\begin{aligned} \text{energy } \int_{B_R} |\nabla \hat{\psi}|^2 &= \int_{\partial B_R} \psi \partial_n \hat{\psi} = \int_{B_R} \nabla \psi \cdot \nabla \hat{\psi} \\ &\Rightarrow \int_{B_R} |\nabla \hat{\psi}|^2 \leq \int_{B_R} |\nabla \psi|^2 \end{aligned}$$

$$\begin{aligned} \text{approx error } \int_{B_R} |\nabla(\psi - \hat{\psi})|^2 &= \int_{\partial B} (\psi - \hat{\psi}) \partial_n \psi \\ &\leq \left(\int_B |\nabla \hat{\psi}|^2 \right)^{1/2} \left(\int_{\partial B} |\nabla \psi|^2 \right)^{1/2} \\ &\Rightarrow \int_{B_R} |\nabla(\psi - \hat{\psi})|^2 \leq \int_{\partial B} |\nabla \psi|^2 \end{aligned}$$

Step 4 Remains to estimate error $\equiv \mathbb{E} \sum_n \frac{\mathbb{1}_{I_n}}{|I_n|} \rho_n^{-d}$.

Preliminary definitions: $\{X_n\}_n$ stationary point process

$$\begin{aligned} &\bullet \mathbb{E} \sum_n h(X_n) = \int h f_1, \quad f_1 = \text{"1-part density"} \\ &\quad \text{or } f_1(x) = \lambda_1 \text{ cst by stationarity} \\ &\quad \text{"intensity" } (\lambda_1 |I_0| = \lambda) \\ &\bullet \mathbb{E} \sum_{n \neq m} h(X_n, X_m) = \iint h f_2, \quad f_2 = \text{"2-part density"} \\ &\quad \& \text{ set } \lambda_2 = \|f_2\|_{\infty} \text{ "2-part intensity"} \\ &\quad \text{e.g. Poisson process: } f_2 = f_1 \otimes f_1 \Rightarrow \lambda_2 = (\lambda_1)^2. \end{aligned}$$

Let's show error $\leq \lambda_2 \log(2 + \lambda_1/\lambda_2)$

Recall $I_n = I(X_n, R_n)$, $\{X_n\}_n \perp \text{cid } \{R_n\}_n$

We compute:

$$\mathbb{E} \sum_n \frac{\mathbb{1}_{I_n}}{|I_n|} \rho_n^{-d} \leq \mathbb{E} \sum_n \mathbb{1}_{|X_n| \leq 1} \underbrace{\rho_n^{-d}}_{\int_{\rho_n}^{\infty} t^{-d-1} dt} \quad (\text{rang } I_n \subset B(X_n))$$

$$\leq \int_1^{\infty} dt \, t^{-d-1} \mathbb{E} \sum_n \mathbb{1}_{|X_n| \leq 1} \underbrace{\mathbb{1}_{\rho_n \leq t}}_{\mathbb{1}_{\exists m \neq n: |X_m - X_n| \leq t}}$$

$$\leq \mathbb{E} \sum_n \mathbb{1}_{|X_n| \leq 1} \leq \lambda_1$$

$$\leq \mathbb{E} \sum_{n \neq m} \mathbb{1}_{|X_n| \leq 1} \mathbb{1}_{|X_n - X_m| \leq t} \leq \lambda_2 t^d$$

$$\leq \int_1^{\infty} dt \, \frac{1}{t} (\lambda_1 t^{-d} + \lambda_2) \leq \lambda_2 \log(2 + \lambda_1/\lambda_2) \quad \square$$